

INTERLAYER AND INTRALAYER SYNCHRONIZATION IN MULTIPLEX FRACTIONAL-ORDER NEURONAL NETWORKS

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Abstract

Fractional-order models describing neuronal dynamics often exhibit better compatibility with diverse neuronal firing patterns that can be observed experimentally. Due to the overarching significance of synchronization in neuronal dynamics, we here study synchronization in multiplex neuronal networks that are composed of fractional-order Hindmarsh–Rose neurons. We compute the average synchronization error numerically for different derivative orders in dependence on the strength of the links within and between network layers. We find that, in general, fractional-order models synchronize better than integer-order models. In particular, we show that the required interlayer and intralayer coupling strengths for interlayer or intralayer synchronization can be weaker if we reduce the derivative order of the model describing the neuronal dynamics. Furthermore, the dependence of the interlayer or intralayer synchronization on the intralayer or interlayer coupling strength vanishes with decreasing derivative order. To support these results analytically, we use the master stability function approach for the considered multiplex fractional-order neuronal networks, by means of which we obtain sufficient conditions for the interlayer and intralayer synchronizations that are in agreement with numerical results.

Keywords: Fractional-Order Neuron Model; Neuronal Network; Multiplex Network; Synchronization; Master Stability Function.

1. INTRODUCTION

Fractional-order differential equations provide a more realistic description of many actual behaviors.¹ Consequently, fractional calculus has been widely used to model diverse dynamical systems.^{2–4} One advantage of the fractional-order modeling is its ability to consider the memory effect

in processes.⁵ Fractional differentiation has been comprehensively considered in different areas like physics and biology.^{6–10} The studies have revealed that the signal propagation in neurons can be well expressed by using fractional-order derivatives.¹¹ Therefore, fractional neuronal models have received substantial attention.^{12–14} For example,

Mondal *et al.*¹⁵ analyzed a FitzHugh–Rinzel neuron model for different fractional derivatives and found diverse firing patterns. Jun *et al.*¹⁶ studied the responses of the fractional Hindmarsh–Rose (HR) system. Their analysis represented that with decreasing the derivative order, the Hopf bifurcation point becomes larger, and the frequency of the firings increases.

One of the most critical problems in the study of the neurons' behaviors is their synchronous firing. Synchronization is a global phenomenon discovered in a variety of dynamic systems.¹⁷ In the brain, synchronization plays significant roles, including information processing and cognitive tasks.¹⁸ Hence, most of the previous researches have been devoted to the synchronous behavior of coupled neurons.^{19–21} Belykh *et al.*²² represented that stable synchronization in bursting neurons is related to their receiving signals. While the synchronization of the linearly coupled neurons only depends on the network topology. Uzuntarla *et al.*²³ studied the synchronization of the bistable neurons under different synaptic connections. They showed that the strong synchronization in the excitatory coupled neurons leads to a spontaneous termination of the neurons' activities. Recently, researchers have also considered the impact of different relating factors on the neurons' synchronization. Some of these factors are the electromagnetic field effects,^{24,25} the time delay,²⁶ the temperature,²⁷ the autapse,²⁸ etc. Among the studies, there are also some works focusing on the synchronization of fractional-order neuron models.^{29–32} Malik and Mir³³ studied the synchronous behavior of two HR systems with fractional derivatives and evaluated their robustness against the noise. Giresse *et al.*³⁴ investigated how the derivatives' order impact on the synchronization of the network of extended HR neurons. They realized that the fractional derivative results in fast synchronization. Meng *et al.*³⁵ considered time-delay fractional-order neurons with electromagnetic induction and derived sufficient conditions for their synchronization analytically.

With the advancement of complex network science, researchers have attended to networks with more complex structures.³⁶ A multilayer network is one of the complex structures that provides a proper representation for many natural systems.³⁷ The multilayer structure allows for considering multiple interactions among systems. Therefore, the synchronizability of multilayer networks has grabbed

considerable attention in different fields such as neuroscience.^{38–40} For instance, Xu *et al.*⁴¹ considered the synchronous behavior of the memristive neurons in a multilayer network. Shafiei *et al.*⁴² focused on synchrony between neurons in a three-layer structure with different synaptic connections. They found opposed synchronization behavior for weak and strong couplings. Rakshit *et al.*⁴³ explored the synchronization of neurons locating in two layers with time-varying intralayer topology. They presented the necessary conditions for the interlayer and intralayer synchronizations using linear stability analysis.

Here, the occurrence of synchronization is under consideration in coupled fractional-order neuron models with multiplex structure. The network is constructed by two layers of small-world structures with identical neurons but non-identical connections. The order of the derivatives and the coupling strengths between and within the layers are the control parameters. The intralayer and the interlayer synchronizations are analyzed by numerical simulations of the network. Then the necessary conditions for synchronizations are derived by extending the master stability function approach to the coupled fractional-order systems in the multiplex framework.

2. THE MODEL

A multiplex network of HR systems with fractional derivatives is considered. The dynamics of model is described by

$$\begin{aligned} D_t^\alpha x &= y + 3x^2 - x^3 - z + I_{\text{ext}}, \\ D_t^\alpha y &= c - 5x^2 - y, \\ D_t^\alpha z &= r(s(x + 1.6) - z), \end{aligned} \quad (1)$$

where D_t^α denotes the α order fractional derivative. The parameters of the model are set at $I_{\text{ext}} = 3.2$, $r = 0.006$, $s = 4$. The firing pattern of model relies on the derivative order α . The bifurcation of the system equation (1) according to the derivative order is demonstrated in Fig. 1a. It can be seen that with a decrement in the derivative order, the peak of the oscillations reduces. Furthermore, the frequency of the spikes in the bursts first increases until $\alpha = 0.934$ and then decreases. Some examples of the time series of the neuron for $\alpha = 1, 0.95, 0.9, 0.85, 0.8, 0.75, 0.7$ are shown in

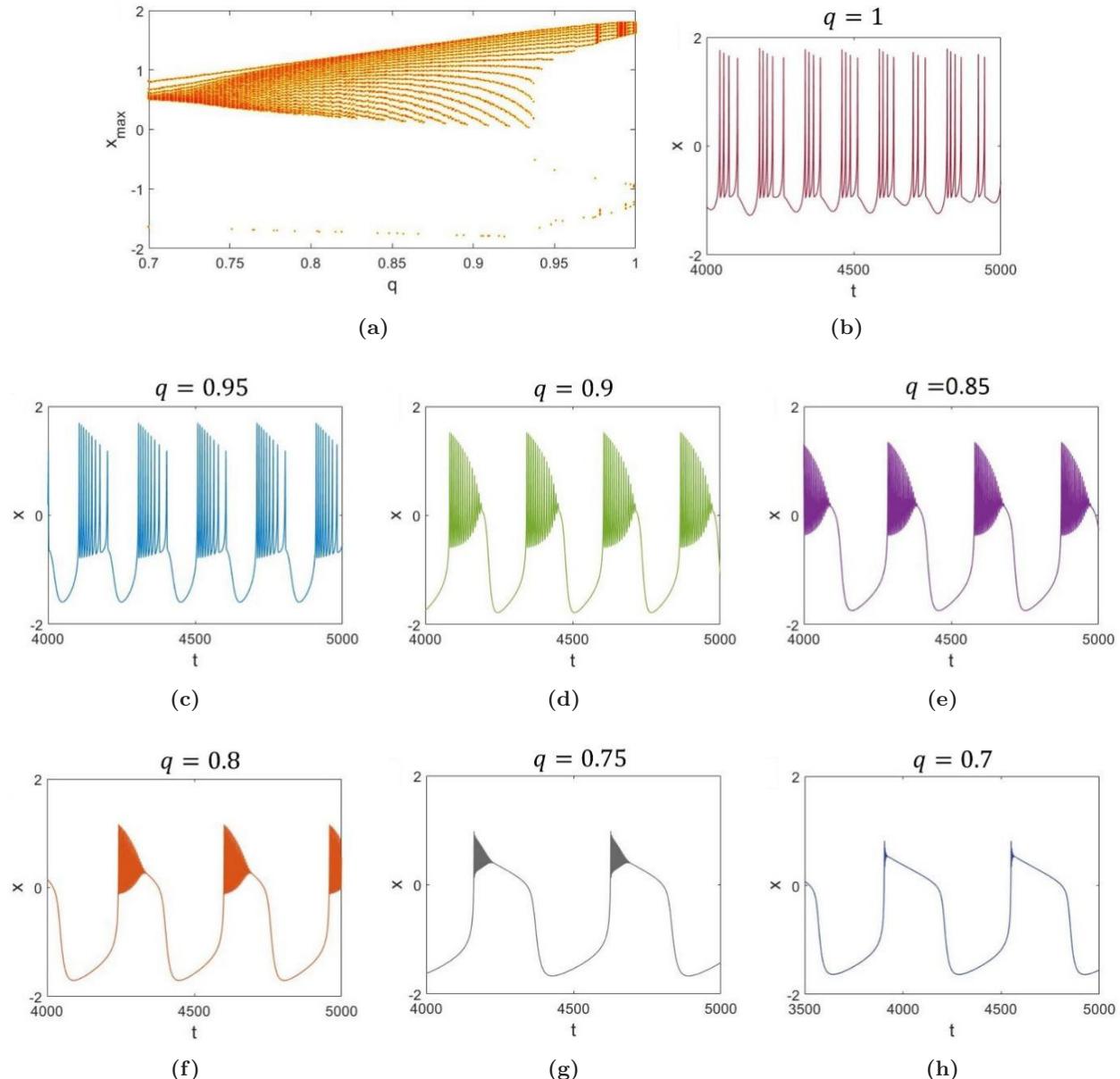


Fig. 1 (a) The bifurcation diagram of the fractional-order HR neuron (Eq. (1)) according to the fractional-order α , (b)–(h) the wave form of the system (Eq. (1)) for different fractional-order derivatives; (b) $\alpha = 1$, (c) $\alpha = 0.95$, (d) $\alpha = 0.9$, (e) $\alpha = 0.85$, (f) $\alpha = 0.8$, (g) $\alpha = 0.75$, (h) $\alpha = 0.7$. As α decreases, the frequency of spikes in bursts first increases and then decreases. While, the frequency of the bursts decreases.

Figs. 1b–1h. It can be seen that as α decreases, the shape of the time series changes from square-wave burster to triangular burster. Moreover, the period of the bursts increases.

$N = 100$ neurons construct each network layer with the small-world structure with probability 0.1 and 10 nearest neighbors' connections. Every node in the first layer is linked with its replica in the second layer. Figure 2 represents a delineative diagram of the network. The equations describing the

network can be given as

$$\begin{aligned} D_t^\alpha X_i &= F(X_i) + \sigma \sum_{j=1}^N G_{ij}^1 h(X_j) \\ &\quad + \epsilon [H(Y_i) - H(X_i)], \\ D_t^\alpha Y_i &= F(Y_i) + \sigma \sum_{j=1}^N G_{ij}^2 h(Y_j) \\ &\quad + \epsilon [H(X_i) - H(Y_i)], \end{aligned} \quad (2)$$

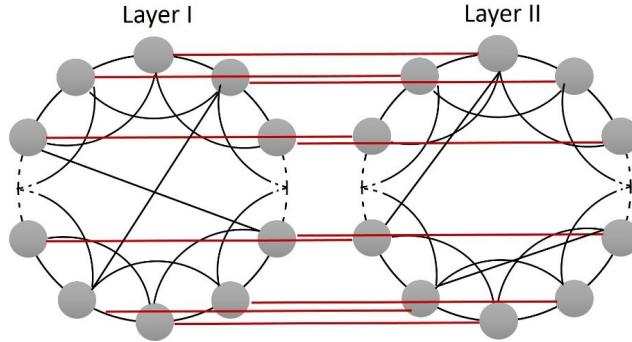


Fig. 2 The schematic of the network that consists of two layers of $N = 100$ neurons with a small-world structure. The black and the red links show the intralayer and the interlayer connections with strength σ and ϵ , respectively.

where X_i and Y_i , $i = 1, \dots, N$ are the m -dimensional state variable of the systems in the first and second layers. D_t^α denotes the α order fractional derivative, $F(X_i)$ is the dynamics of the single HR system as Eq. (1), σ and ϵ are the intralayer and interlayer coupling strengths. The connection matrixes are denoted by $G^1 = [G_{ij}^1]_{N \times N}$ and $G^2 = [G_{ij}^2]_{N \times N}$, which are zero-row-sum matrixes with $G_{ij} = 1$ if the nodes i and j are linked, and $G_{ij} = 0$, else. The inner coupling functions are shown by $h(X)$ and $H(X)$, which are assumed to be linear through x variables, thus $h(X) = H(X) = \begin{bmatrix} x & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$. The paper aims to find the effects of the strength of the connections (σ, ϵ) and the derivative order (α) on the synchrony of the neurons.

3. RESULTS

The two-layer network is solved for different derivative orders, and the emergence of synchronization within and between the layers is considered by varying the coupling strengths. The fourth-order Runge–Kutta is applied to solve integer-order equations. The fractional-order model is considered based on the Caputo fractional derivative operator. For a given fractional differential equation as

$$\begin{aligned} {}_0^C D_t^\alpha x(t) &= F(t, x(t)), \\ x(0) &= x_0, \end{aligned} \quad (3)$$

${}_0^C D_t^\alpha$ expresses the α Caputo fractional order as

$$\begin{aligned} {}_0^C D_t^\alpha x(t) &= \frac{1}{\Gamma(n-\alpha)} \int_0^t f^n(\tau)(t-\tau)^{n-\alpha-1} d\tau, \\ n-1 < \alpha < n, \quad n \in N. \end{aligned} \quad (4)$$

For the numerical solving of the fractional-order equations the Adams–Bashforth–Moulton method described in Ref. 44 is used.

In the following sections, the appearance of intralayer and interlayer synchronizations in the network is investigated.

3.1. Intralayer Synchronization

Intralayer synchronization occurs when the neurons within the layers become synchronous. To evaluate the intralayer synchronization numerically, the averaged intralayer synchronization error is computed as follows:

$$E_X = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \sum_{j=2}^N \frac{\|X_j(t) - X_i(t)\|}{N-1} dt, \quad (5)$$

where $\|\cdot\|$ represents the Euclidean norm, and T is a long duration of time. Figure 3 represents the intralayer synchronization error for different derivative orders. Figure 3a shows that the intralayer synchronization in the integer-order model depends on both coupling strengths (σ and ϵ). It can be seen that when ϵ increases from 0.1 to 0.4, the intralayer error decays to zero for lower σ . But when ϵ reaches 0.5, the synchrony threshold (zero error point) becomes larger. Thus, by increasing ϵ , the required intralayer coupling strength (σ) for intralayer synchronization first decreases and then increases. By changing the derivatives to the fractional order, the trend of the error changes. For $\alpha = 0.9$ (Fig. 3a), the intralayer synchronization is obtained for lower σ than the integer-order case. Furthermore, the dependence on the interlayer coupling strength decreases such that the zero points of error are almost the same for all ϵ values. Figures 3c and 3d show that the dependence of intralayer synchrony on ϵ decreases for lower orders. Moreover, for $\alpha = 0.8$, the intralayer synchrony is achieved for slightly lower σ than $\alpha = 0.9$. While for $\alpha = 0.7$, more σ is needed for synchrony.

To confirm the results of the synchronization error, the necessary conditions for the synchronization of the network are also derived by applying the master stability function method. We consider the synchronization within the layers, regardless of the synchronization between two layers. Assuming the synchronous manifold of the first layer and the second layer to be $S_X(t)$ and $S_Y(t)$, thus, $X_i(t) = S_X(t)$ and $Y_i(t) = S_Y(t)$ for $i = 1, \dots, N$,

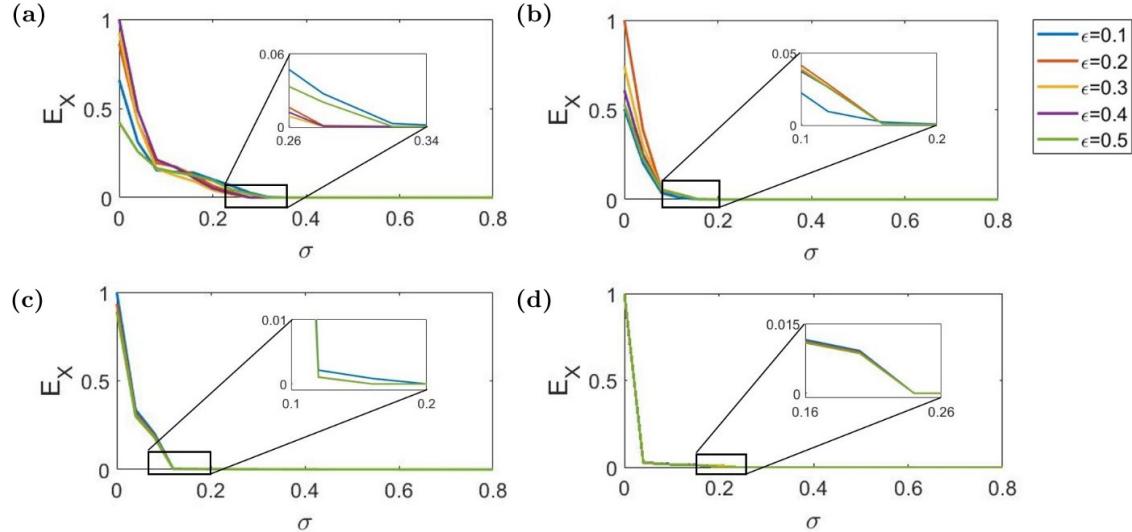


Fig. 3 The intralayer synchronization error according to intralayer coupling strength (σ), for different interlayer coupling strengths (blue: $\epsilon = 0.1$, red: $\epsilon = 0.2$, yellow: $\epsilon = 0.3$, purple: $\epsilon = 0.4$, green: $\epsilon = 0.5$). (a) $\alpha = 1$; (b) $\alpha = 0.9$; (c) $\alpha = 0.8$; (d) $\alpha = 0.7$. With decreasing α , the intralayer synchronization is achieved in weaker couplings.

where

$$\begin{aligned} D_t^\alpha S_X &= F(S_X) + \epsilon[H(S_Y) - H(S_X)], \\ D_t^\alpha S_Y &= F(S_Y) + \epsilon[H(S_X) - H(S_Y)]. \end{aligned} \quad (6)$$

Then the perturbations from the synchronous manifold can be defined as $\delta X_i = X_i - S_X$ and $\delta Y_i = Y_i - S_Y$. By derivation of the perturbed equations, we get $D_t^\alpha \delta X_i = D_t^\alpha X_i - D_t^\alpha S_X$ and $D_t^\alpha \delta Y_i = D_t^\alpha Y_i - D_t^\alpha S_Y$. Thus, the linearized equations can

be obtained as follows:

$$\begin{aligned} D_t^\alpha \delta X_i &= JF(S_X) \delta X_i + \sigma \sum_{j=1}^N G_{ij}^1 Jh(S_X) \delta X_j \\ &\quad + \epsilon [JH(S_Y) \delta Y_i - JH(S_X) \delta X_i], \\ D_t^\alpha \delta Y_i &= JF(S_Y) \delta Y_i + \sigma \sum_{j=1}^N G_{ij}^2 Jh(S_Y) \delta Y_j \\ &\quad + \epsilon [JH(S_X) \delta X_i - JH(S_Y) \delta Y_i]. \end{aligned} \quad (7)$$

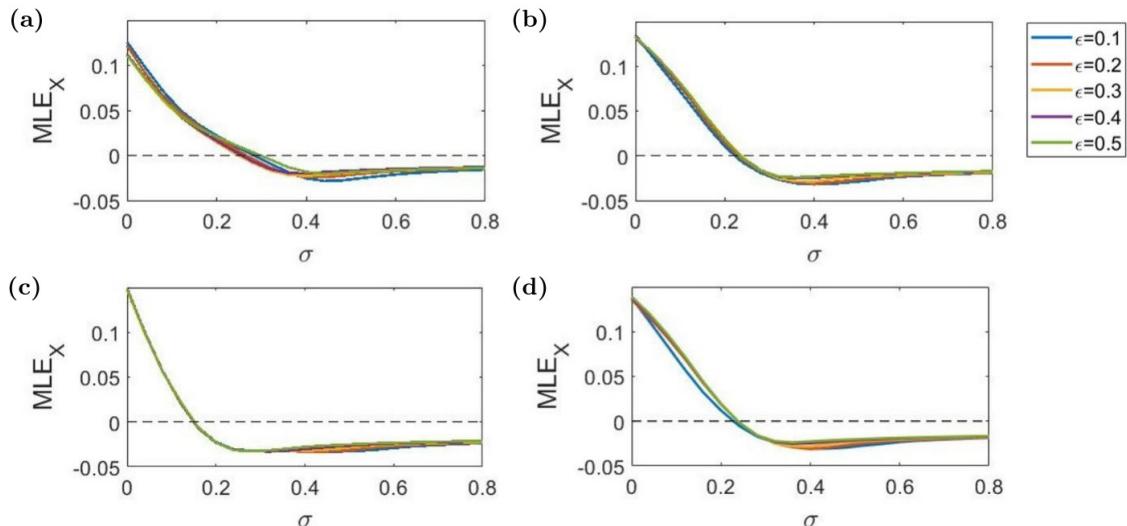


Fig. 4 The MLE of the linearized equations obtained for the intralayer synchronization (Eq. (7)) according to intralayer coupling strength (σ), for different interlayer coupling strengths (blue: $\epsilon = 0.1$, red: $\epsilon = 0.2$, yellow: $\epsilon = 0.3$, purple: $\epsilon = 0.4$, green: $\epsilon = 0.5$). (a) $\alpha = 1$; (b) $\alpha = 0.9$; (c) $\alpha = 0.8$; (d) $\alpha = 0.7$.

The maximum Lyapunov exponent (MLE) of the linearized system (Eq. (7)), which is transverse to the synchronous manifold (S_X and S_Y) specifies the stability of the intralayer synchrony. When $\text{MLE} < 0$, the linearized perturbation system is stable about zero equilibrium point and thus, the intralayer synchronization is stable.

Figure 4 shows the MLE of Eq. (7) by altering the intralayer coupling strength (σ) for different derivative orders. It is observed that by reducing the fractional order, the required σ for intralayer synchronization decreases until $\alpha = 0.8$ and then increases for $\alpha = 0.7$. The dependence of the MLE on ϵ also reduces with decreasing α . These results are well matched with the ones obtained from the numerical intralayer synchronization error.

3.2. Interlayer Synchronization

The interlayer synchronization is known as the synchrony between the corresponding nodes in two layers, regardless of the synchrony within layers.

The interlayer synchronization error is found by

$$E_{XY} = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \sum_{j=1}^N \frac{\|Y_j(t) - X_j(t)\|}{N} dt, \quad (8)$$

with $\|\cdot\|$ being the Euclidean norm and T as a long time interval. The interlayer synchronization error for $\alpha = 1, 0.9$ and 0.8 is demonstrated in Fig. 5. Figure 5a shows that for $\sigma = 0.05$, the interlayer synchronization of the integer network is achieved for $\epsilon > 0.32$. By increasing the intralayer coupling strength, the interlayer synchrony is obtained for stronger ϵ . For example, for $\sigma = 0.2$ and 0.5 , the network shows the interlayer synchronization for $\epsilon > 0.4$ and 0.48 , respectively. When α changes to 0.9 , no considerable change happens for the synchronization threshold for $\sigma = 0.05$ and 0.2 . But the threshold for the stronger intralayer coupling strengths is influenced. It can be seen in Fig. 5b that the interlayer synchrony is obtained for $\epsilon > 0.4$ when $\sigma = 0.5$. Therefore, the dependence

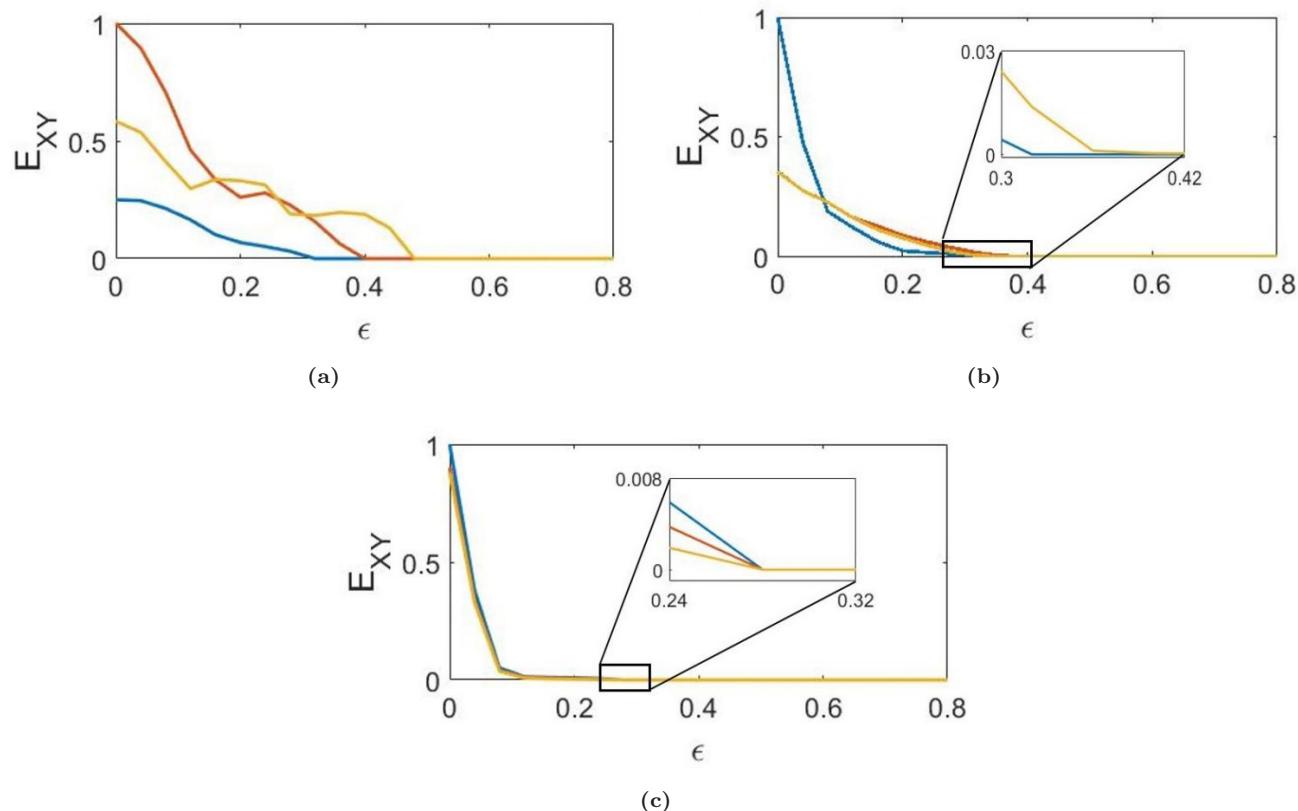


Fig. 5 The interlayer synchronization error according to interlayer coupling strength (ϵ), for different intralayer coupling strengths (blue: $\sigma = 0.05$, red: $\sigma = 0.2$, yellow: $\sigma = 0.5$). (a) $\alpha = 1$; (b) $\alpha = 0.9$; (c) $\alpha = 0.8$. The interlayer synchronization threshold is lessened by decreasing α .

of interlayer synchrony threshold on the intralayer coupling strength reduces for high σ values. With decrement of α to 0.8, the interlayer synchrony does not depend on the σ value and emerges for $\epsilon > 0.28$.

Next, we apply the master stability function for the synchronization between two layers. The synchronous manifold in this case is $S = X_i = Y_i$, which is obtained from the following equation:

$$D_t^\alpha S = F(S) + \sigma \sum_{j=1}^N G_{ij}^1 h(S). \quad (9)$$

Thus, the perturbation can be defined as $\eta_i = Y_i - X_i$. The derivative of the perturbation is $D_t^\alpha \eta_i = D_t^\alpha Y_i - D_t^\alpha X_i$, which leads to

$$\begin{aligned} D_t^\alpha \eta_i &= F(Y_i) - F(X_i) - 2\epsilon H(\eta_i) \\ &\quad + \sigma \sum_{j=1}^N G_{ij}^2 h(Y_j) - \sigma \sum_{j=1}^N G_{ij}^1 h(X_j). \end{aligned} \quad (10)$$

Considering $\Delta G_{ij} = G_{ij}^2 - G_{ij}^1$, we have

$$\begin{aligned} D_t^\alpha \eta_i &= F(Y_i) - F(X_i) - 2\epsilon H(\eta_i) \\ &\quad + \sigma \sum_{j=1}^N G_{ij}^2 [h(Y_j) - h(X_j)] \\ &\quad - \sigma \sum_{j=1}^N \Delta G_{ij} h(X_j), \end{aligned} \quad (11)$$

which leads to the following linearized equation:

$$\begin{aligned} D_t^\alpha \eta_i &= [JF(S_i) - 2\epsilon JH(S_i)]\eta_i \\ &\quad + \sigma \sum_{j=1}^N G_{ij}^2 Jh(S_j)\eta_j - \sigma \sum_{j=1}^N \Delta G_{ij} h(S_j). \end{aligned} \quad (12)$$

Since there is a slight difference between the topologies of the layers, we can neglect the term $\sigma \sum_{j=1}^N \Delta G_{ij} h(S_j)$ in the interlayer synchronization. Consequently, the sign of the MLE of Eq. (12), transverse to the synchronous manifold Eq. (9),

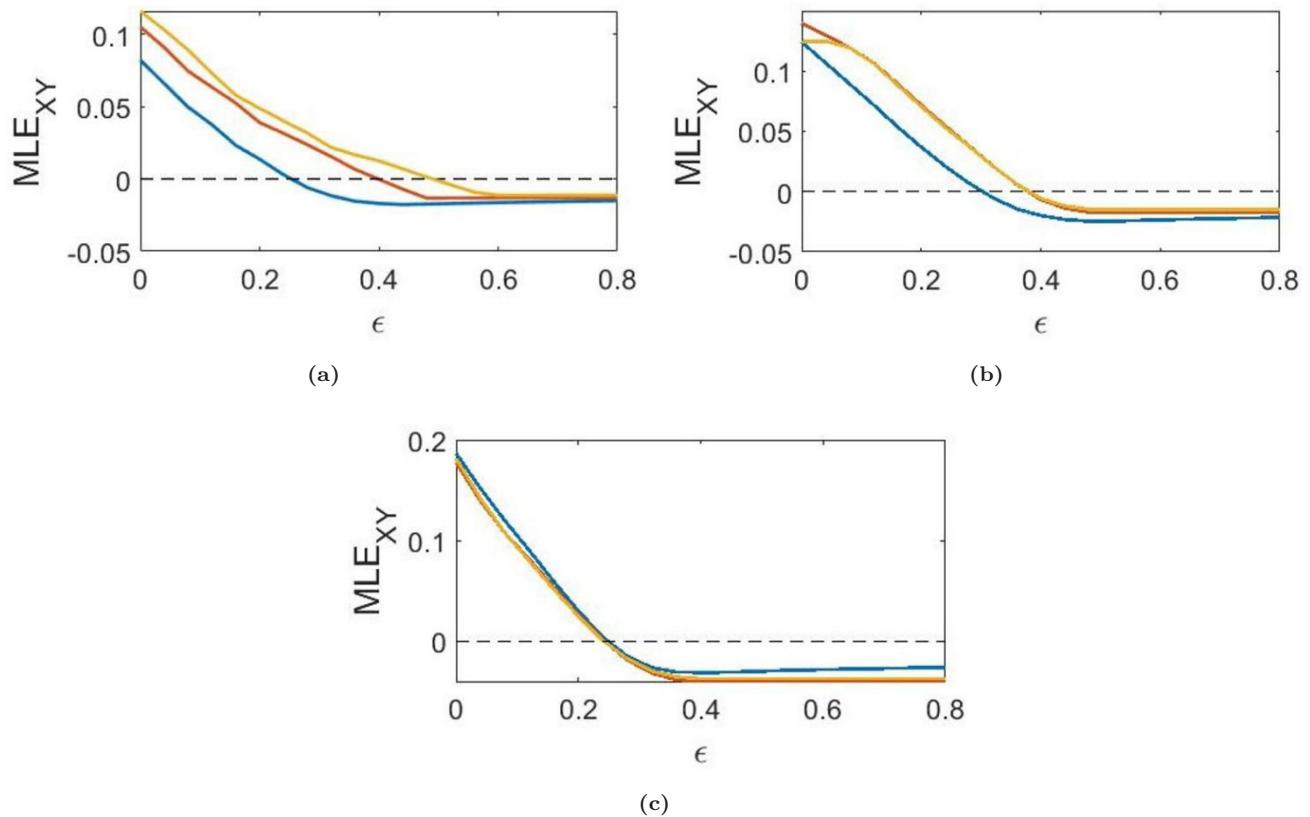


Fig. 6 The MLE of the linearized equations obtained for the interlayer synchronization (Eq. (12)) according to interlayer coupling strength (ϵ), for different intralayer coupling strengths (blue: $\sigma = 0.05$, red: $\sigma = 0.2$, yellow: $\sigma = 0.5$). **(a)** $\alpha = 1$; **(b)** $\alpha = 0.9$; **(c)** $\alpha = 0.8$.

determines the stability of the interlayer synchronization.

Figure 6 represents the MLE of Eq. (12) according to interlayer coupling strength. The figure demonstrates that the required ϵ for interlayer synchronization decreases with decreasing the derivative order. Furthermore, the dependence of the interlayer synchrony threshold on the intralayer coupling strength disappears.

4. CONCLUSION

This paper studied a multiplex network of HR models with fractional derivatives from the viewpoint of interlayer and intralayer synchronization. The small-world structure constructed each layer of the network, and both the interlayer and intralayer couplings were considered diffusive. The integer and fractional networks with different orders were solved numerically, and the averaged intralayer, and interlayer synchronization errors were calculated for different coupling strengths. Moreover, the master stability function was extended to obtain the necessary conditions for the interlayer and intralayer synchronizations in the multiplex fractional network. The results showed that for the integer network, the intralayer synchronization threshold first reduces and then increases with decreasing the interlayer coupling strength. In addition, the interlayer synchronization is achieved for stronger interlayer couplings as the intralayer coupling strength is enlarged. In comparison with the integer network, in the fractional network, both types of synchronizations were acquired in weaker couplings. By decreasing the order of the derivatives, the synchronization appeared within the layers in lower intralayer coupling strengths and between the layers in lower interlayer coupling strengths. Furthermore, the necessary interlayer (or intralayer) coupling strength for the interlayer (or intralayer) synchronization became independent of the intralayer (or interlayer) coupling strength.

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